

# **Simulation of diffusion processes with discontinuous coefficients**

Antoine Lejay

Projet TOSCA, INRIA Nancy Grand-Est, Institut Élie Cartan

From collaborations with Pierre Étoré  
and Miguel Martinez

## Divergence form operators in modelling

In many problems, a concentration, pressure, size of a population, ... is given by a partial differential equation (PDE) that follows from

$$\int_V (C(t + \Delta t, x) - C(t, x)) dx = \Delta t \int_{\partial V} \mathbf{J}(t, x) d\sigma(x)$$

with

- $V$  volume
- $\mathbf{J}(t, x)$  flux
- $C(t, x)$  concentration

Then

$$\frac{\partial C(t, x)}{\partial t} = \text{div}(\mathbf{J}(t, x))$$

In general, the flux can be related to the concentration itself by

$$\mathbf{J}(t, x) = a(x)\nabla C(t, x)$$

for a matrix  $a(x)$  (diffusivity/permeability/...), which leads to

$$\frac{\partial C(t, x)}{\partial t} = \text{div}(a(x)\nabla C(t, x)) \text{ or } \text{div}(a(x)\nabla C(x)) = f(x) \text{ (steady state)}$$

Under the assumption (**uniform ellipticity**) that for some constants  $\lambda, \Lambda > 0$ ,

$$\lambda|\xi|^2 \leq a(x)\xi \cdot \xi \leq \Lambda|\xi|^2, \quad \forall \xi \in \mathbb{R}^d, \quad \forall x \in \mathbb{R}^d$$

and  $a$  is measurable (no continuity assumption !) the PDE

$$\begin{cases} \frac{\partial u(t, x)}{\partial t} = \frac{1}{2} \operatorname{div}(a(x)\nabla u(t, x)) \\ u(0, x) = g(x) \end{cases}$$

has a unique solution (in the weak sense), that is

$$\begin{aligned} \int_{\mathbb{R}^d} \phi(0, x)g(x) dx + \int_0^T \int_{\mathbb{R}^d} \frac{\partial \phi(t, x)}{\partial t} u(t, x) dt \\ = \frac{1}{2} \int_0^T \int_{\mathbb{R}^d} a(x)\nabla u(t, x)\nabla \phi(t, x) dx dt \end{aligned}$$

for all  $\phi \in \mathcal{C}^{\infty, \infty}([0, T]; \mathbb{R}^d)$ ,  $\phi(T, x) = 0$ .

This solution is  $(\alpha/2, \alpha)$ -Hölder continuous and weakly differentiable, but one cannot expect **better regularity** in general.

However, if  $a$  is locally regular, then  $u$  is also be locally regular.

Consider the case where  $S$  is a surface smooth enough, and  $a$  is smooth on each sides  $S_+$  and  $S_-$  of  $S$ .

Then, for  $x$  on  $S$ ,

$$u(t, x+) = u(t, x-)$$

$$\text{and } \underbrace{a(x+)\mathbf{n}_+ \cdot \nabla u(t, x+) = a(x-)\mathbf{n}_- \cdot \nabla u(t, x-)}_{\text{continuity of the flux}}$$

Examples of models with discontinuities/interfaces:

- Concentration of a fluid in a porous media (Darcy's law) with different type of rocks
- Diffusion of species in several type of habitats
- Composite materials
- ...

## Divergence form op. and diffusion processes

There exists a **fundamental solution**  $p(t, x, y)$  for the differential operator  $\frac{1}{2} \operatorname{div}(a \nabla \cdot)$  so that the solution to

$$\begin{cases} \frac{\partial u(t, x)}{\partial t} = \frac{1}{2} \operatorname{div}(a(x) \nabla u(t, x)) \\ u(0, x) = g(x) \end{cases}$$

may be written

$$u(t, x) = \int_{\mathbb{R}^d} p(t, x, y) g(y) dy.$$

In addition,  $0 \leq p(t, x, y)$ ,  $\int_{\mathbb{R}^d} p(t, x, y) dy = 1$  and

$$p(t, x, y) \leq \frac{C_1}{t^{d/2}} \exp\left(-\frac{C_2 |y - x|^2}{2t}\right) \quad (\text{Nash-Aronson estimate})$$

$\implies \frac{1}{2} \operatorname{div}(a \nabla \cdot)$  is the infinitesimal generator of a strong Markov, continuous stochastic process  $X$ .

## Link with solutions of SDEs

If  $a$  is of class  $\mathcal{C}^1$ , then

$$L = \frac{1}{2} \operatorname{div}(a \nabla \cdot) = \frac{1}{2} \sum_{i,j=1}^d a_{i,j} \frac{\partial^2}{\partial x_i \partial x_j} + \sum_{i,j=1}^d \frac{1}{2} \frac{\partial a_{i,j}}{\partial x_i} \frac{\partial}{\partial x_j}$$

and then  $X$  is solution to the SDE

$$X_t = x + \int_0^t \sigma(X_s) dB_s + \int_0^t \frac{1}{2} \nabla a(X_s) ds$$

with  $\sigma \sigma^T = a$  and  $B$  is a Brownian motion.

In general,  $X$  is **not** a semi-martingale, because  $\nabla a$  has no meaning.

However, the differential operator acts locally, so that in the region where  $a$  is smooth,  $X$  behaves like a “good” diffusion.

What happens at the surface of discontinuity?

## Divergence and non-divergence form operators

$$D = \frac{1}{2} \operatorname{div}(a \nabla \cdot)$$

- There always exists a process associated to  $D$ , as soon as  $a$  is measurable, uniformly elliptic bounded.
- Gaussian estimates are the keys
- Difficult to simulate
- Martingale + something
- Natural class of functions: Sobolev spaces

$$N = \frac{1}{2} \sum_{i,j=1}^d a_{i,j} \partial_{i,j}^2$$

- Existence and uniqueness is provided under regularity of  $a$  (Hölder continuity) or on restrictive conditions on the discontinuities.
- Uniqueness may fail with discontinuous coefficients
- Easy to simulate
- Semi-Martingales
- Natural class of functions:  $\mathcal{C}^2$

**Note:** The special case of  $N = \frac{1}{2} a_{i,j} \partial_{i,j}^2$  with  $a(x) = \rho(x) \operatorname{Id}$ ,  $\rho : \mathbb{R}^d \rightarrow \mathbb{R}$  can be understood as a special case of divergence form operator with invariant measure  $\rho$  and then regularity assumptions on  $\rho$  may be dropped.

## A simple case: $d = 1$

From now, we work under the simple case of  $d = 1$ .

☹ It is too **restrictive** to understand what happens when  $d > 1$ , which is the case for most of real-world problems.

☺ But it allows one to have a **better understanding** on what happens.

From now, we consider

$$L = \frac{\rho}{2} \nabla (a \nabla \cdot)$$

where

- $\rho, a$  are measurable
- $\lambda \leq \rho(x) \leq \Lambda$  for all  $x \in \mathbb{R}$
- $\lambda \leq a(x) \leq \Lambda$  for all  $x \in \mathbb{R}$
- $a$  and  $\rho$  are piecewise smooth and have a left/right limit at any point

Our choice of  $L$  covers both the case of divergence and non-divergence form operators



## Remark 1: the drift is not a problem

If one wishes to consider

$$L = \frac{\rho}{2} \nabla(a \nabla \cdot) + b \nabla$$

with a measurable, bounded  $b$ , one has only to remark that

$$L = \frac{\rho e^{-\Phi}}{2} \nabla(a e^{\Phi} \nabla \cdot)$$

with

$$\Phi(x) = 2 \int^x \frac{b(y)}{a(y)\rho(y)} dy$$

and  $a e^{\Phi}$  and  $\rho e^{-\Phi}$  satisfy locally the previous hypotheses.

Thus, it is not a problem to consider a drift, or to set  $b = 0$ .

## Remark 2: on the transmission condition

A function  $f$  in the domain  $\text{Dom}(L)$  of  $L$  satisfies

$$f(x-) = f(x+) \text{ and } a(x-)\nabla f(x-) = a(x+)\nabla f(x+)$$

at any point  $x$  where  $a$  or  $\rho$  is discontinuous.

If  $a$  and  $\rho$  are continuous on  $(-\infty, \beta)$  and on  $(\beta, +\infty)$ , set

$$\tilde{a}(x) := \begin{cases} \lambda a(x) & \text{on } (-\infty, \beta) \\ a(x) & \text{otherwise} \end{cases} \text{ and } \tilde{\rho}(x) := \begin{cases} \rho(x)/\lambda & \text{on } (-\infty, \beta) \\ \rho(x) & \text{otherwise} \end{cases}$$

so that

$$\tilde{L} = \frac{\tilde{\rho}}{2} \nabla(\tilde{a} \nabla)$$

is such that

$$\tilde{L}f(x) = Lf(x) \text{ for } x \in (-\infty, \beta) \cup (\beta, +\infty)$$

but for  $f$  in  $\text{Dom}(\tilde{L})$ ,

$$f(x-) = f(x+) \text{ and } \lambda a(x-)\nabla f(x-) = a(x+)\nabla f(x+)$$

## On one-dimensional diffusion processes

The operator  $L$  is the infinitesimal generator of a strong Markov process with continuous paths.

This process can be constructed either as the process associated to the divergence-form operator  $L$ , or by its scale function and its speed measure (or by other means).

**Scale function:** There exists a continuous, increasing function  $S$  (unique up to additive and multiplicative constants) such that for  $x < y < z$

$$\mathbb{P}_y [\tau_x < \tau_z] = \frac{S(z) - S(y)}{S(z) - S(x)}$$

**Speed measure:** There exists a measure  $\mu$  such that

$$\mathbb{E}_y [\tau_{[x,z]}] = \int_x^z G_{x,z}(y, u) \mu(du)$$

for the Green function  $G_{x,y}(y, u)$  on  $[-x, z]$  of  $L$  (given by a simple expression).

## On one-dimensional diffusion processes

The scale function is related to the probability that the diffusion reaches a given point  $x$  before another given point  $y$ , while the speed measure gives an indication to the average times it takes to exit from some interval.

If  $B$  is a Brownian motion, the processes  $t \mapsto B_t$  and  $t \mapsto B_{2t}$  have the same scale function  $S(x) = x$  but different speed measures.

In dimension one, any diffusion process  $X$  is **fully characterized** by its scale function  $S$  and its speed measure  $\mu$

There exists a random time change  $T$  related to  $m$  such that

$$X_t = S(B_{T(t)}).$$

## Convergence results

If  $a^n$ ,  $b^n$  and  $\rho^n$  are such that

$$\frac{1}{a^n} \xrightarrow[n \rightarrow \infty]{L^2_{\text{loc}}(\mathbb{R})} \frac{1}{a}, \quad \frac{1}{\rho^n} \xrightarrow[n \rightarrow \infty]{L^2_{\text{loc}}(\mathbb{R})} \frac{1}{\rho} \quad \text{and} \quad \frac{b^n}{a^n \rho^n} \xrightarrow[n \rightarrow \infty]{L^2_{\text{loc}}(\mathbb{R})} \frac{b}{a\rho}$$

then the process  $X^n$  associated to  $L^n = \frac{\rho^n}{2} \nabla(a^n \nabla \cdot) + b^n \nabla$  converges in distribution to the process  $X$  associated to  $L = \frac{\rho}{2} \nabla(a \nabla \cdot) + b \nabla$  for any starting point.

A way to deal with the problem of discontinuous coefficients consists in using a sequence of mollifiers to regularize the coefficients, but

☹ From the numerical point of view, it does not work very well

☹ From the theoretical point of view, there is nothing to understand

Instead, we will assume that  $a$  and  $\rho$  are **piecewise constant**, which can be a good approximation of piecewise smooth function by adding a lot of small jumps.

## Local behavior of the process

If  $a$  and  $\rho$  are constant on the intervals  $(x_i, x_{i+1})$  then for  $f \in \text{Dom}(L)$ ,

$$\begin{cases} Lf(x) = \frac{1}{2}\rho(x_i+)a(x_i+)f''(x) \text{ on } (x_i, x_{i+1}), \\ f \in \mathcal{C}^2((x_i, x_{i+1})), \quad \forall i \\ f(x_i-) = f(x_i+), \\ (1 - q_i)\nabla f(x_i-) = (1 + q_i)\nabla f(x_i+) \text{ with } q_i = \frac{a(x_i+) - a(x_i-)}{a(x_i+) + a(x_i-)}. \end{cases}$$

With the remark on the transmission condition above (multiply  $a$  by  $\lambda$  on  $(x_i, x_{i+1})$  and  $\rho$  by  $1/\lambda$  on  $(x_i, x_{i+1})$ ), the problem consists in finding the local behavior of a process  $X$  such that

- It is a Brownian motion with a given speed on  $(x_{i-1}, x_i)$
- It is a Brownian motion with a given speed on  $(x_i, x_{i+1})$
- The functions in the domain of its infinitesimal generator satisfy  $(1 - q_i)\nabla f(x_i-) = (1 + q_i)\nabla f(x_i+)$  and  $f(x_i-) = f(x_i+)$  at  $x_i$

## A simple case

**Assumption:**  $a$  and  $\rho$  are constant on  $\mathbb{R}_-^*$  and  $\mathbb{R}_+^*$ .

It can be showned (by several means) that

$$X_t = x + \int_0^t \sqrt{a(X_s)\rho(X_s)} dB_s + \underbrace{\frac{a(0+) - a(0-)}{a(0+) + a(0-)}}_{\in(-1,1)} L_t^0(X),$$

where

- $B$  is a Brownian motion
- $L_t^0(X)$  is the symmetric **local time** at 0 of  $X$ :

$$L_t^0(X) = \lim_{\epsilon \rightarrow 0} \frac{1}{2\epsilon} \int_0^t \mathbf{1}_{[-\epsilon, \epsilon]}(X_s) ds$$

The local time characterizes the time spend by  $X$  at 0.

$t \mapsto L_t^0(X)$  is continuous and non-decreasing. However, it increases only on  $\{t \geq 0 \mid X_t = 0\}$  which has a Lebesgue measure equal to 0!

The local time has an effect only when  $X$  reaches 0. Otherwise,  $X$  behaves like a Brownian motion with diffusion coefficient  $a\rho/2$ .

## A simple case: why such an SDE?

**Heuristic:** consider a function of class  $\mathcal{C}_b^2(\mathbb{R} \setminus \{0\}) \cap \mathcal{C}(\mathbb{R})$  with  $a(0+)\nabla f(0+) = a(0-)\nabla f(0-)$

Itô-Tanaka formula  $\implies$

$$f(X_t) = f(x) + \underbrace{\int_0^t \frac{1}{2} (f'(X_s+) + f'(X_s-)) dX_s}_{\int_0^t f'(X_s) \sqrt{\sigma(X_s)a(X_s)} dB_s + K} + \underbrace{\frac{1}{2} \int_0^t (f'(X_s+) - f'(X_s-)) dL_t^0(X_s)}_{=-K} + \frac{1}{2} \int_0^t \sigma(X_s) a(X_s) f''(X_s) ds$$

so that  $M_t = f(X_t) - f(x) - \int_0^t \frac{1}{2} \int_0^t \sigma(X_s) a(X_s) f''(X_s) ds$  is a martingale with brackets  $\langle M \rangle_t = \int_0^t a(X_s) \sigma(X_s) f'(X_s)^2 ds$

$\implies$  characterization of the infinitesimal generator of  $X$ .



## SDE with local time

Conditions for strong existence and convergence results for SDE of type

$$X_t = x + \int_0^t \sigma(X_s) dB_s + \int_{\mathbb{R}} \nu(dy) L_t^y(X)$$

have been provided to J.-F. Le Gall in the 80's.

In our case, we consider measures of type

$$\nu(dx) = g(x) dx + \sum_i \alpha_i \delta_{x_i}$$

With the occupation density formula,

$$\int_{\mathbb{R}} g(x) L_t^x(X) dx = \int_0^t \sigma(X_s)^2 g(X_s) ds$$

so that one can consider measures of type  $\nu(dx) = \sum_i \alpha_i \delta_{x_i}$ .

This class of SDE is stable under application of one-to-one mappings  $\Phi$

such that  $\Phi \in \mathcal{C}(\mathbb{R}) \cap \mathcal{C}^2(\mathbb{R} \setminus \{z_i\})$  with

$$(1 + q_i) \nabla \Phi(z_i+) = (1 - q_i) \nabla \Phi(z_i-).$$

## The Skew Brownian motion

The **Skew Brownian motion** is a generalization of the Brownian motion that depends on a parameter  $\alpha \in [-1, 1]$ .

$$a = \begin{cases} a_+ & \text{on } \mathbb{R}_+ \\ a_- & \text{on } \mathbb{R}_-, \end{cases}, \quad \rho = \begin{cases} 1/a_+ & \text{on } \mathbb{R}_+ \\ 1/a_- & \text{on } \mathbb{R}_- \end{cases} \quad \text{and } \alpha = \frac{a_+}{a_+ + a_-}$$


The process  $X$  is solution to the SDE (Harrison-Shepp)


$$X_t = x + B_t + \beta L_t^0(X) \quad \text{with } \beta = 2\alpha - 1.$$

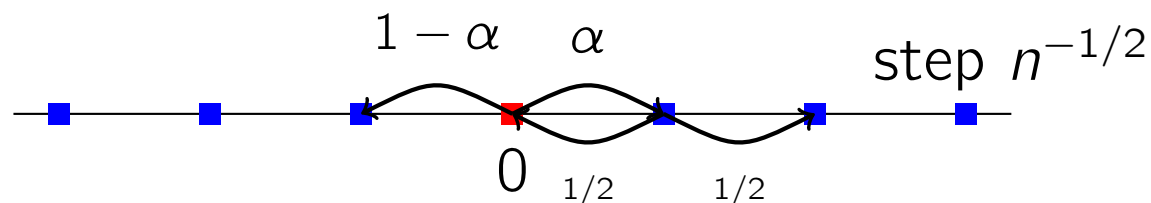
 A possible **construction** of the Skew Brownian motion (Itô-McKean)

- Consider the excursions of a reflected Brownian motion
- Change the sign of each excursion with independent Bernouilli random variables of parameter  $\alpha$

## Constructions of the Skew Brownian motion

 Another possible **construction** of the Skew Brownian motion (N. Portenko, 78') it is the process whose infinitesimal generator is
 
$$L = \frac{1}{2}\Delta + a\delta_0$$

 Yet another **construction** of the Skew Brownian motion (Harrison-Shepp) consider a simple random walk  $S_n$  on  $\mathbb{Z}$  such that
 
$$\mathbb{P}[S_{n+1} = 1 | S_n = 0] = \alpha.$$
 Then  $n^{-1/2}S_{nt} \xrightarrow[n \rightarrow \infty]{\text{dist.}} \text{SBM}(\alpha)$



 We have also a simple expression for the density of the SBM.

There are indeed many ways to construct a SBM, and some schemes follows easily from these constructions.

The Skew Brownian motion is the **basic element** to understand SDE with local time.

## An apparent paradox?

What can be done in the case of  $L = \frac{\rho}{2} \nabla(a \nabla \cdot)$ ?

Heuristically, the discontinuity may be interpreted as **permeable barrier**: the particle goes to one side or the other with a given probability.



However, the situation is more complex.

Consider

$$a(x) = \begin{cases} a_+ & \text{if } x \geq 0, \\ a_- & \text{if } x < 0, \end{cases} \text{ and } \rho = 1.$$

Then for any  $h > 0$ ,

$$\mathbb{P}_0 [\tau_h < \tau_{-h}] = \frac{a_+}{a_+ + a_-}$$

Yet for any  $t > 0$ ,

$$\mathbb{P}_0 [X_t > 0] = \frac{\sqrt{a_+}}{\sqrt{a_+} + \sqrt{a_-}}$$

## Simulation techniques

In order to get simulation techniques for

$$L = \frac{\rho}{2} \nabla (a \nabla \cdot)$$

a possible approach consists in using well-chosen one-to-one mapping to change the process.

### Suppression of the local time

$Y_t = S(X_t)$  with  $S'(x) = 1/a(x)$  is solution to

$$Y_t = \int_0^t \sqrt{\frac{\rho \circ S^{-1}(Y_s)}{a \circ S^{-1}(Y_s)}} dB_s$$

**Simulation:** Euler scheme for SDE with discontinuous coefficients

**i**L. Yan, *Ann. Appl. Probab.* 2002 for the convergence of the Euler scheme with discontinuous coeff.

**i**M. Martinez, Ph.D. thesis, 2004 for application to divergence form operators and computation of the rate of convergence.

## Reduction to the SBM

$$\Psi(x) = \int^x \frac{1}{\sqrt{\rho(y)a(y)}} dy$$

then  $Y_t = \Psi(X_t)$  is solution to

$$Y_t = \Phi(x) + B_t + \underbrace{\frac{\sqrt{a(0+)/\rho(0+)} - \sqrt{a(0-)/\rho(0-)}}{\sqrt{a(0+)/\rho(0+)} + \sqrt{a(0-)/\rho(0-)}}}_{2\alpha-1} L_t^{\Psi(0)}(Y)$$

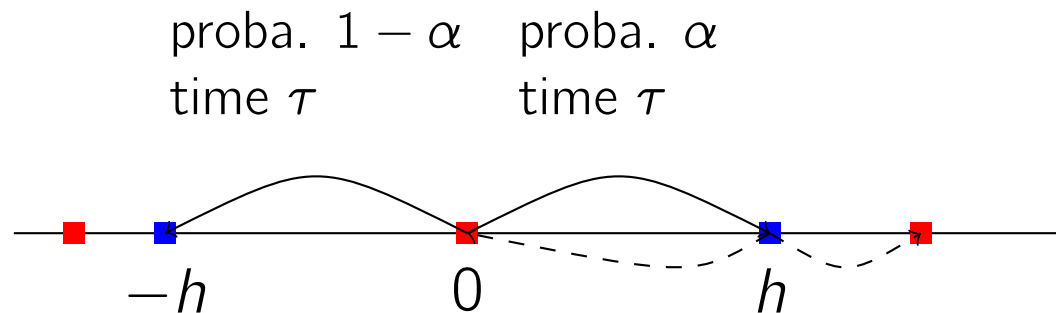
(To simplify, we assume only one point of discontinuity at 0)

The process  $Y$  is then (locally if there are more than one point of discontinuity) a skew Brownian motion.

If at time  $t$  the process is at 0 and

$$\tau = \inf \{ s > t \mid X_s = \pm h \}$$

then  $\tau$  and  $Y_\tau$  are independent,  $\mathbb{P} [ Y_\tau = h ] = \alpha$  and  $\tau$  is distributed as the first exit time of the Brownian motion from  $[-h, h]$



**Simulation:** At 0 (or any point of discontinuity), one can re-inject the particle in  $-h$  or  $h$  by simulating  $Y_\tau$  and increasing the time by  $\tau$ .

It is also possible to simulate exactly  $(\tau \wedge T, Y_{\tau \wedge T})$  (and to do so for the Brownian motion for any starting point in a given interval).

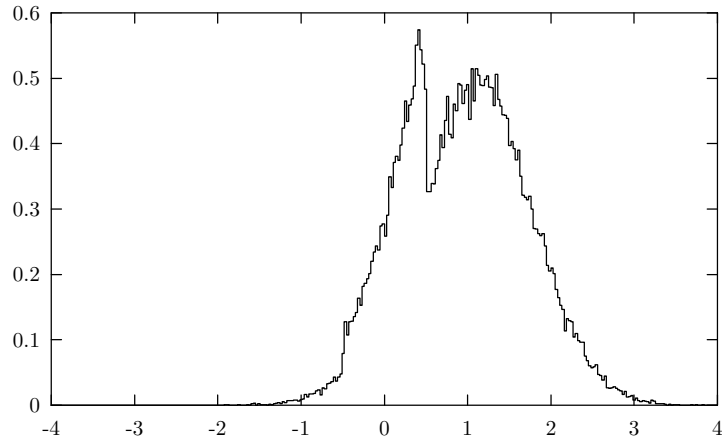
This method becomes time consuming unless the coefficients are piecewise constant with not too many points of discontinuities.

**i** A.L. & M. Martinez, *Ann. Appl. Probab.*, 2005.

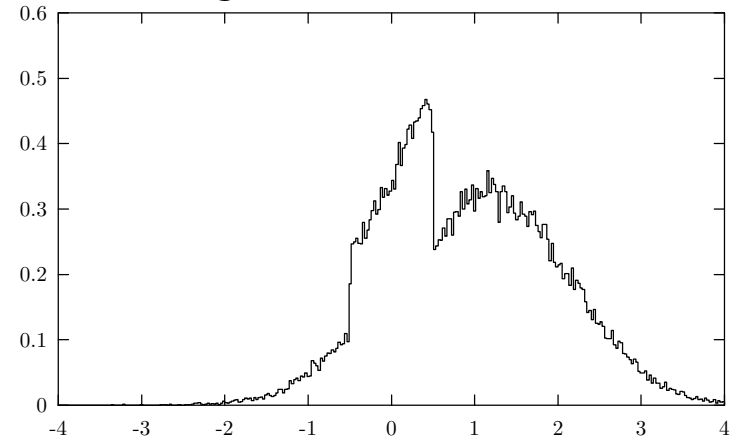


## A numerical example: doubly SBM

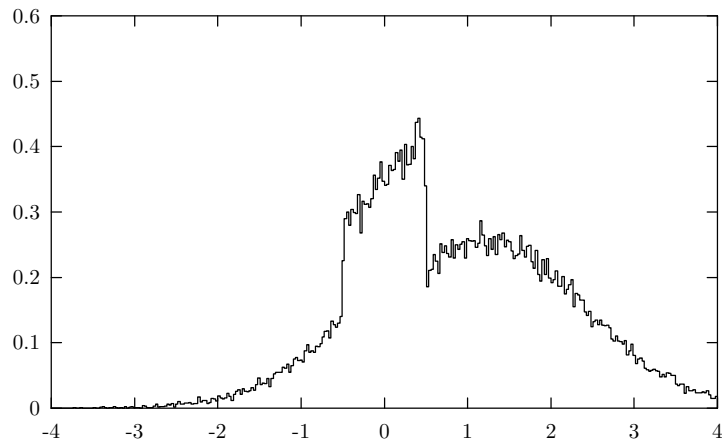
$$X_t = 1 + B_t + \frac{2}{3}L_t^{-1/2}(X) - \frac{2}{3}L_t^{1/2}(X)$$



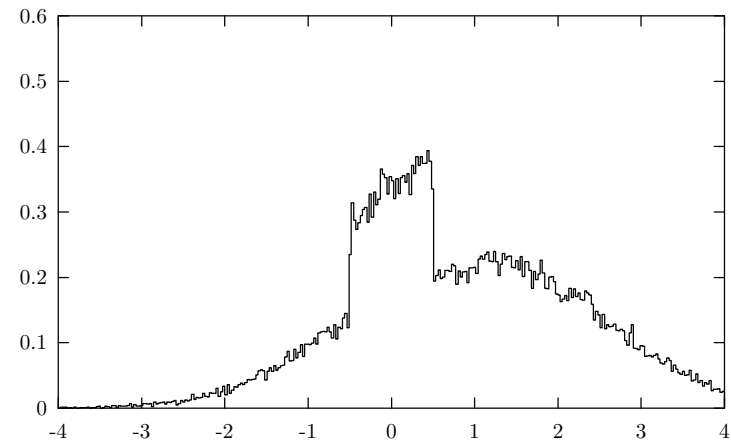
$t = 1/2$



$t = 1$



$t = 3/2$



$t = 2$

## Approximation by a random walk I

If  $a$  and  $\rho$  are piecewise constant, up to a small perturbation of the positions of the points of discontinuities, it is possible to assume that there are contained  $\{x_i\}_{i \in \mathbb{Z}}$  with  $\Psi(x_i) = ih$  for a given step  $h$ , where

$$\Psi(x) = \int_0^x \frac{1}{\sqrt{\rho(y)a(y)}} dy.$$

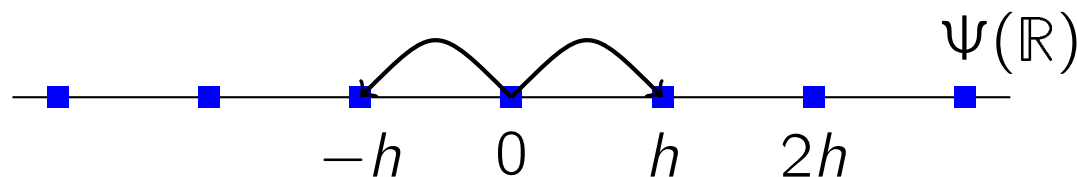
$Y := \Psi(X)$  is solution to

$$Y_t = \Psi(x) + B_t + \sum_{i \in \mathbb{Z}} \beta_i L_t^{ih}(Y).$$

where

$$\beta_i = \frac{\sqrt{a(x_i+)/\rho(x_i+)} - \sqrt{a(x_i-)/\rho(x_i-)}}{\sqrt{a(x_i+)/\rho(x_i+)} + \sqrt{a(x_i-)/\rho(x_i-)}}$$

$1 - \beta_0$     $\beta_0$   
time  $h^2$    time  $h^2$



**Simulation:** walk on the grid  $\{ih\}_{i \in \mathbb{Z}}$

- Perform a random walk  $(\xi_k)$  with

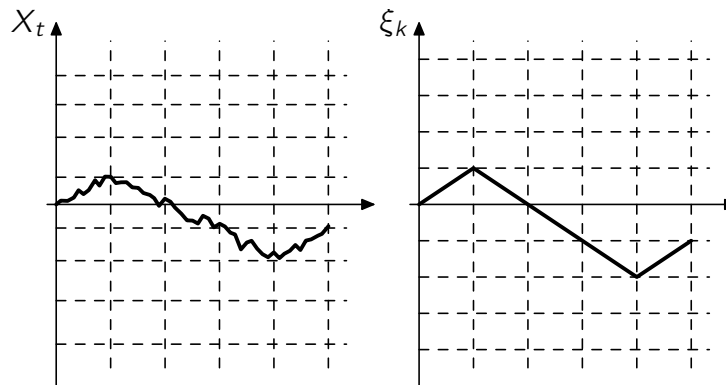
$$\mathbb{P}[\xi_{k+1} = (i+1)h \mid \xi_k = ih] = \frac{1 + \beta_i}{2}$$

- At each step, increase the time by

$$h^2 = \mathbb{E}[\tau_i \mid \xi_k = ih, \xi_{k+1} = (i+1)h] = \mathbb{E}[\tau_i \mid \xi_k = ih, \xi_{k+1} = (i-1)h]$$

with  $\tau_i = \inf \{ t > 0 \mid |Y_t - \xi_i| = h \}$

- ☞ The position  $\Psi^{-1}(\xi_k)$  represents an approximation of the position of  $X_t$  at time  $t = kh^2$ .



Rate of convergence as in Donsker's theorem

📄 P. Étoré, *Electron. J. Probab.*, 2006

## Approximation by a random walk II

The difficulty with the previous method is that the grid depends on the coefficients.

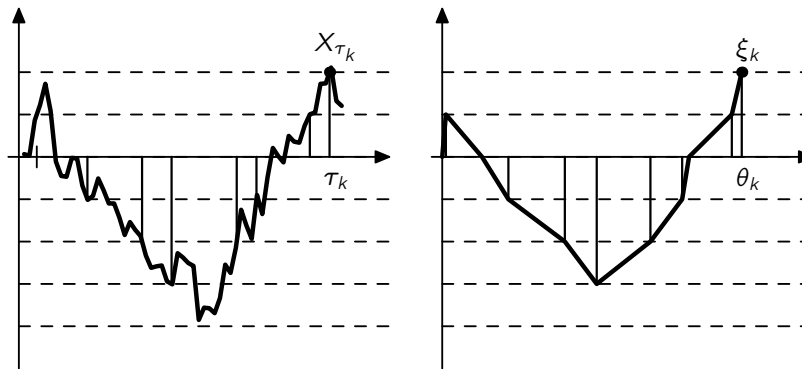
Consider a process  $X$  with infinitesimal generator  $L = \frac{\rho}{2} \nabla (a \nabla \cdot)$  whose coefficients are uniformly elliptic and bounded (this method does not use SDE with local times)

Consider a grid  $\mathcal{G} = \{x_i\}_i$  with  $x_{i-1} < x_i$  for all  $i$ .

Set

$$\delta_i = \inf \{ t > 0 \mid X_t \in \mathcal{G} \setminus \{X_{\tau_i}\} \}, \quad \tau_{i+1} = \tau_i + \delta_i$$

The  $\tau_i$ 's represent the time at which  $X$  reaches successive levels.



## Simulation: random walk on the grid $\mathcal{G}$

- Perform a random walk with

$$\mathbb{P} [\xi_{k+1} = x_{i+1} \mid \xi_k = x_i] = \mathbb{P}_{x_i} [X_{\tau_k} = x_{i+1}] = u(x_i)$$

with

$$\begin{cases} Lu = 0 \text{ on } (x_{i-1}, x_{i+1}), \\ u(x_{i+1}) = 1, \quad u(x_{i-1}) = 0. \end{cases}$$

- At each step, increment the time by

$$t_k := \mathbb{E} [\tau_k \mid \xi_k = x_i, \xi_{k+1}] = v(x_i)/w(x_i) \text{ with}$$

$$\begin{cases} Lv = -w & \text{on } (x_{i-1}, x_{i+1}) \\ v(x_{i-1}) = v(x_{i+1}) = 0 \end{cases}$$

with

$$w = \begin{cases} u & \text{if } \xi_i = x_k, \quad \xi_{i+1} = x_{k+1}, \\ 1 - u & \text{if } \xi_i = x_k, \quad \xi_{i+1} = x_{k-1}. \end{cases}$$

☞ At time  $\theta_i = \sum_{j < i} t_j$ ,  $(\theta_i, \xi_i)$  is an approximation of  $(\tau_i, X_{\tau_i})$ .

📖 P. Étoré & A.L., *ESAIM Probab. Stat.*, 2007.

## Other approaches and related works

- ★ J. M. Ramirez, E.A. Thomann & E.C. Waymire (Oregon State University) used the multi-skew Brownian motion in a geophysical context.
- ★ M. Decamps, A. De Scheeper, M. Goovaerts & W. Schoutens (Catholic University of Leuven) have proposed some “self-exciting threshold interest rates models” in finance that relies on the SBM and numerical methods using perturbation formulas of the density of the SBM.
- ★ N. Limić (University of Zagreb) have proposed a scheme for  $d > 1$  where the infinitesimal generator of a continuous time Markov process is constructed from a finite-volume scheme.